Problem 5

An energy difference between the spin-singlet (S=0) state and the spintriplet (S=1) state is given by

 $E(\uparrow\downarrow\downarrow\downarrow\downarrow\uparrow)-E(\uparrow\uparrow,\leftarrow\leftarrow,\downarrow\downarrow)=2J$

on the basis of the atomic orbital method. Here, 2J = -4St + 2J'.

 $\mathbf{J'} = \iint \phi_{\mathfrak{s}}(1)\phi_{\mathfrak{s}}(1)\phi_{\mathfrak{s}}(2)\phi_{\mathfrak{s}}(2)\Big($

and <u>S</u> is the overlap $+\frac{e^2}{r_{12}}$) $dv_i dv_i$ integral.

This energy difference is described by the following effective Hamiltonian using the respective spin-operators S_a and S_b at the A and B sites.

 $H_{\rm eff} = -2J \, s_{\rm a} \cdot s_{\rm b}$

Noting that $S = s_a + s_b$ and $S(S+1)=S^2 = s_a^2 + s_b^2 + 2 s_a \cdot s_b$, when S=0 and S=1, show that $s_a \cdot s_b = -3/4$ and $s_a \cdot s_b = 1/4$,

respectively, and the energy difference between S=0 and S=1 is given by 2*J*. The above Hamiltonian is known as the Heisenberg Hamiltonian is written by $H= -2J\Sigma_{ii}S_i \cdot S_i$

(6-1)The trial wave functions of Molecular Orbital and Heitler-London (HL) with spin-singlet and HL with spin-triplet are given by

$$\begin{split} \Phi_{s} &= \psi_{1}(1)\psi_{1}(2) \\ \Phi_{\text{HL}} &= \frac{\psi_{a}(1)\psi_{b}(2) + \psi_{b}(1)\psi_{a}(2)}{\sqrt{2}} \end{split} \\ \Phi_{T} &= \frac{\psi_{a}(1)\psi_{b}(2) - \psi_{b}(1)\psi_{a}(2)}{\sqrt{2}} \end{split}$$

Show the respective expected eigen energies are given as follows;

$$E_{s} = 2(\varepsilon + t) + \frac{1}{2}U_{0} + \frac{1}{2}U_{1} + J'$$
$$E_{\text{HL}} = 2\varepsilon + U_{1} + J'$$
$$E_{T} = 2\varepsilon + U_{1} - J'$$

Problem 6

We next consider a more realistic solution of the electronic state in H_2 beyond either Heitler – London or Molecular Orbital approximations, by assuming H(1,2)=h(1)+h(2)+V(1,2) as a Hamiltonian for H_2 where h(1)+h(2) is the sum of one-electron Hamiltonian and V(1,2) is the Coulomb repulsive interaction between electrons. Several integrals are defined as the follows;

$$\varepsilon \equiv \int \phi_a(1)h(1)\phi_a(1)dv_1 = \int \phi_b(1)h(1)\phi_b(1)dv_1$$
$$t \equiv \int \phi_a(1)h(1)\phi_b(1)dv_1 = \int \phi_b(1)h(1)\phi_a(1)dv_1$$
$$U_0 \equiv \int \phi_a(1)^2\phi_a(2)^2V(1,2)dv_1dv_2 = \int \phi_b(1)^1\phi_b(2)^2V(1,2)dv_1dv_2$$

 $U_1 \equiv \int \phi_a(1)^2 \phi_b(2)^2 V(1, 2) dv_1 dv_2 = \int \phi_b(1)^2 \phi_a(2)^2 V(1, 2) dv_1 dv_2$

 $J' \equiv \int \phi_a(1)\phi_b(1)\phi_a(2)\phi_b(2)V(1,2)dv_1dv_2$

Solve the next problems.

(6-2) In order to improve the HL wave function;

$$\varPhi_{\rm HL} = \frac{1}{\sqrt{2(1+S^2)}} [\phi_a(1)\phi_b(2) + \phi_b(1)\phi_a(2)]$$

We incorporate the following states, using $\phi_a(1) \phi_a(2)$ and $\phi_b(1) \phi_b(2)$

$$\Phi' = \frac{1}{\sqrt{2}} [\psi_a(1)\psi_a(2) + \psi_b(1)\psi_b(2)]$$

Then, the improved trial wave function is expressed by $\boldsymbol{\varphi} = c_1 \boldsymbol{\varphi}_{HL} + c_2 \boldsymbol{\varphi}'$. Show how to get the following relation and solve an eigen energy for this state

$$\begin{vmatrix} 2\varepsilon + U_1 + J' - E & 2t \\ 2t & 2\varepsilon + U_0 + J' - E \end{vmatrix} = 0$$

If $U_0 - U_1 > /t / is$ valid for Mott insulator, show that the eigen energy is given as the follow;

$$E_{\text{HL}} = 2\varepsilon + U_1 + J' - \frac{4t^2}{U_0 - U_1}$$